

A Photoelectrochemical Study of Polyacetylene, (CH)_x

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Synopsis. The photoelectrochemical behavior of *p*-type semiconducting *trans*-(CH)_x film in the presence of *N,N'*-dimethyl-4,4'-bipyridinium as the solution species is described. The spectral response of the photocurrent shows a mismatch with the absorption spectrum, and the quantum yield ($\approx 10^{-3}$) of the charge flow is extremely low. It is assumed that (CH)_x film has a high trap density and that the lifetime of the photogenerated hole is limited by trapping and recombination in the space-charge region.

In recent years, there has been a considerable interest in the electrical and optical properties of polyacetylene, (CH)_x, film as a new class of semiconducting material.^{1–6} A (CH)_x-sodium polysulfide solution photovoltaic cell has been fabricated as an active photoelectrode for a photoelectrochemical cell.⁷ Although Chen *et al.* mentioned that the (CH)_x photocathode exhibited significant photoresponse with an open circuit voltage, $V_{oc} \approx 0.3$ V, and a short circuit current, $I_{sc} \approx 40 \mu\text{A}/\text{cm}^2$, under illumination of approximately 1 sun,⁷ the photoresponse ($V_{oc} \approx 60$ mV, $I_{sc} \approx 1 \mu\text{A}/\text{cm}^2$) in our work was very low. These results suggest that, in order to use the (CH)_x as an active photoelectrode for a photoelectrochemical cell, it is of primary importance to characterize its general photoelectrochemical behavior before any modification experiments are performed. The studies reported here concern the basic photoelectrochemical properties of the *trans*-(CH)_x film in the presence of *N,N'*-dimethyl-4,4'-bipyridinium (MV²⁺) as the solution species.

Experimental

The preparation of the *trans*-(CH)_x film has been described elsewhere.⁸ The film was about 0.1 mm thick. The conductivity of the film was found to be about $10^{-6} \Omega^{-1} \text{cm}^{-1}$, as determined by the standard four-probe van der Pauw technique.⁹ The ohmic contact with a thin copper sheet was obtained by the use of Electrodag on the shiny side of the (CH)_x film. The preparation of the (CH)_x working electrode was done according to Chen's method.⁷

Photocurrent measurements were performed under potentiostatic conditions with a home-made potentiostat.

The light source used in the study of the photoelectrochemical effect was a 100-W high-pressure mercury lamp with a glass filter ($\lambda \leq 430$ nm cut-off). A water filter with a 10-cm optical pathlength as a heat-absorbing filter was employed. Chemical actinometry for 436-nm-wavelength light, obtained by combination with a KL-43 filter, was carried out using the potassium ferrioxalate system. The measurement of the action spectrum of the photocurrent was done by the lock-in (NF Model LI-574) technique using a modulation of the light beam with 8 Hz. As a light source in this case, a Xe 500-W lamp was used, while a grating monochromator (Nikon G-250) was employed for the wavelength selection.

Reagent-grade chemicals were used without further purification. All the solutions were deoxygenated for at least

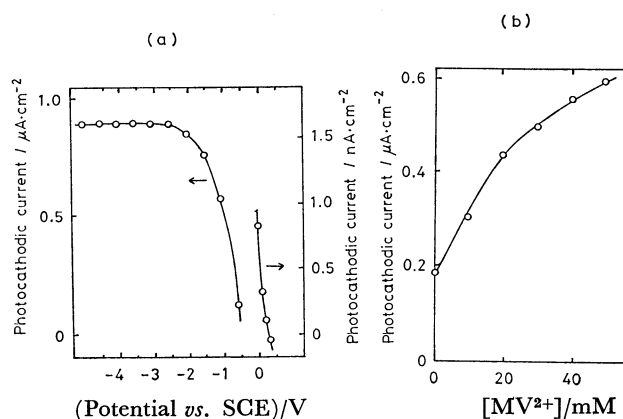


Fig. 1. (a) Steady-state photocurrent vs. electrode potential for 20 mM MV²⁺ in 1 M KCl (pH 5) at *trans*-(CH)_x electrode; (b) Steady-state photocurrent as a function of concentration of MV²⁺ (1 M KCl, pH 5, $U_{SCE} = -1.0$ V).

30 min with purified nitrogen before each experiment. All the experiments were carried out with the solution under nitrogen without stirring.

Results and Discussion

Cyclic-voltammetric curve revealed that the (CH)_x film did not react with MV²⁺ in the dark at all. Figures 1(a) and (b) show the steady-state photocurrent-voltage characteristic and the dependence of the signal on the MV²⁺ concentration respectively. The results can be qualitatively explained in terms of the *p*-type semiconducting properties of the (CH)_x film. The cathodic photocurrent indicates that MV²⁺ acts as a primary acceptor of the electron as the minority carrier ejected from the (CH)_x into the solution and can be attributed to the reduction of MV²⁺ to the purple MV^{•+}, as evidenced by the blue-purple color streaming from the (CH)_x film surface at $U_{SCE} \leq -1.5$ V. This fact suggests that the lower edge of the (CH)_x conduction band is positioned at energies above the standard redox potential (-0.7 V vs. SCE) of MV²⁺/MV^{•+}.¹⁰ Figure 1(b) shows the observation of the signal in the absence of MV²⁺ as well. The electron acceptor in this case remains unidentified, but dissolved oxygen gas as an impurity is a possible candidate, because O₂ can mediate the transfer of the conduction-band electron across the (CH)_x film-solution interface. The onset potential ($U_{SCE} \approx 0.4$ V) of the photocurrent, which approximately corresponded to the flat-band potential for the (CH)_x electrode, was almost independent of the solution pH (1–13), and the signal decreased with the solution acidity when U_{SCE} was kept constant. When a solution containing sodium polysulfide (20 mM as Na₂S₂) and 1 M (mol dm⁻³) KCl at pH 11.6 was used,⁷ the photo-

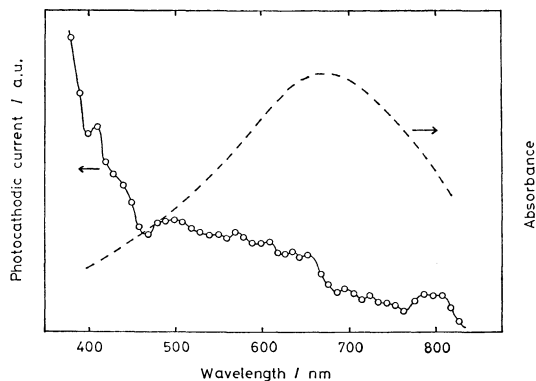


Fig. 2. Spectral response (in arbitrary units) of *trans*-(CH)_x electrode in 20 mM MV²⁺ + 1 M KCl aqueous solution (pH 5).

Dashed curve indicates absorption spectrum of *trans*-(CH)_x with about 1 μm in thickness.

current was lower ($\approx 0.27 \mu\text{A}/\text{cm}^2$ at $U_{\text{SCE}} = -1.0 \text{ V}$) than that of the 20 mM MV²⁺ system at pH 5 ($\approx 0.45 \mu\text{A}/\text{cm}^2$ at $U_{\text{SCE}} = -1.0 \text{ V}$). This lower photocurrent may be mainly caused by an absorbance of the polysulfide solution used, and we may conclude that, in the polysulfide electrolyte, (CH)_x does not exhibit the significant photoresponse previously reported.⁷⁾

In Figure 2 the spectral response of the (CH)_x electrode is shown. The data on the vertical axis were obtained by dividing the photocurrent at a certain wavelength by the number of photons incident on the electrode. It can be seen that the spectral response of the photocurrent shows a mismatch with the absorption spectrum of the (CH)_x film, which exhibits an absorption maximum at 670 nm.¹¹⁾ Similar behavior has been observed for the solid-state (CH)_x Schottky junction.¹²⁾ The threshold of the photocurrent response is observed at about 830 nm ($\approx 1.5 \text{ eV}$), in agreement with the direct-band gap calculated from the optical studies of the (CH)_x film.¹⁾ The anticorrelation between the action and the absorption spectrum indicates that the light absorbed on the front surface of the (CH)_x film is not effective in creating free-charge carriers. In effect, most holes generated on the front surface are not long-lived enough to diffuse into the space-charge layer, while holes generated within the bulk have a finite probability of transiting the space-charge layer to be registered as carriers in the external circuit. The rate of charge generation at a certain distance, x , from the front surface is proportional to $\exp(-kx)$, where k is the absorption coefficient of the (CH)_x film (for example, $k_{670} \approx 2 \times 10^5 \text{ cm}^{-1}$).¹¹⁾ Detailed studies of the photovoltaic response for *p*-(CH)_x: *n*-CdS heterojunction have implied the existence of a meta-stable trapping state 0.9 eV below the conduction band of the (CH)_x film.⁶⁾ Therefore, the anticorrelation may be explained by assuming that the (CH)_x film has an excessively high trap density; if the penetration depth of the light ($\lambda \leq 600 \text{ nm}$) is large, the trapping of the photogenerated holes may be insignificant compared to the case of about a 670 nm-wavelength light corresponding to the absorption maximum of the (CH)_x film. The quantum yield of the charge flow under

the 436-nm-light irradiation was extremely low, being 1×10^{-3} at $U_{\text{SCE}} = -1.0 \text{ V}$ in the 20 mM MV²⁺ + 1 M KCl aqueous solution (pH 5). Furthermore, the light-intensity dependence of the photocurrent varied with the applied potential; the light-intensity exponent for the photocurrent was less than unity and increased from 0.8 to 0.9 as $-U_{\text{SCE}}$ increased from 0.5 to 3.0 V. These results may be related to the recombination or trapping of the photogenerated holes in the bulk of the (CH)_x film, which arises from the high trap density and which well limits the photocurrent to a significant level, if we consider that an increase in the electric field in the space-charge layer would cause the photogenerated holes to have an increasing chance of transiting the space-charge layer before being trapped.¹³⁾

References

- 1) C. K. Chiang, C. R. Fincher, Jr., Y. W. Park, A. J. Heeger, H. Shirakawa, E. J. Louis, S. C. Gau, and A. G. MacDiarmid, *Phys. Rev. Lett.*, **39**, 1098 (1977).
- 2) C. K. Chiang, Y. W. Park, A. J. Heeger, H. Shirakawa, E. J. Louis, and A. G. MacDiarmid, *J. Chem. Phys.*, **69**, 5098 (1978).
- 3) C. R. Fincher, Jr., M. Ozaki, M. Tanaka, D. Peebles, L. Lauchlan, A. J. Heeger, and A. G. MacDiarmid, *Phys. Rev. B*, **20**, 1589 (1979).
- 4) Y. W. Park, A. Denesteyn, C. K. Chiang, A. J. Heeger, and A. G. MacDiarmid, *Solid State Commun.*, **29**, 747 (1979).
- 5) M. Ozaki, D. L. Peebles, B. R. Weinberger, C. K. Chiang, S. C. Gau, A. J. Heeger, and A. G. MacDiarmid, *Appl. Phys. Lett.*, **35**, 83 (1979).
- 6) M. Ozaki, D. L. Peebles, B. R. Weinberger, A. J. Heeger, and A. G. MacDiarmid, *J. Appl. Phys.*, **51**, 4252 (1980).
- 7) S. N. Chen, A. J. Heeger, Z. Kiss, A. G. MacDiarmid, S. C. Gau, and D. L. Peebles, *Appl. Phys. Lett.*, **36**, 96 (1980).
- 8) T. Ito, H. Shirakawa, and S. Ikeda, *J. Polym. Sci., Polym. Chem. Ed.*, **12**, 11 (1974).
- 9) L. J. van der Pauw, *Philips Tech. Rev.*, **20**, 220 (1958); J. Lange, *J. Appl. Phys.*, **35**, 2659 (1964).
- 10) S. Hunig, J. Gross, and W. Schenk, *Justus Liebigs Ann. Chem.*, **1**, 324 (1973).
- 11) H. Shirakawa, T. Ito, and S. Ikeda, *Polym. J.*, **4**, 460 (1973).
- 12) T. Tani, P. M. Grant, W. D. Gill, G. B. Street, and T. C. Clarke, *Solid State Commun.*, **33**, 499 (1980).
- 13) The AsF₆-doped *trans*-(CH)_x film with $8.4 \times 10^{-3} \Omega^{-1} \text{ cm}^{-1}$ exhibited 5×10^{16} and $2.5 \times 10^{18} \text{ cm}^{-3}$ as the ionized acceptor concentration (N_A) and the trap density (N_T), respectively.¹⁴⁾ Accordingly, it may be reasonable to assume $N_A \approx 10^{15} \text{ cm}^{-3}$ for the present (CH)_x film with $10^{-6} \Omega^{-1} \text{ cm}^{-1}$. Then, about 10^4 \AA as the space-charge layer width under the band bending of about 1 eV can be calculated, when $\epsilon \approx 10$ is taken as the dielectric constant. In this case, most of the absorption maximum-photons ($k_{670} \approx 2 \times 10^5 \text{ cm}^{-1}$) incident on the (CH)_x film would be absorbed within 10^3 \AA of the front surface, and most of the photogenerated holes would be not long-lived enough to transit the space-charge layer with a 10^4 \AA width, due to the trapping by the high-trap density (probably $N_T > 2.5 \times 10^{18} \text{ cm}^{-3}$). This may lead to the absence of any parallel effect for the photocurrent signal, as is shown in Fig. 2.
- 14) H. Shirakawa and T. Tani, *Kotai Butsuri*, **14**, 435 (1979).